

# Applications and Investigations into Complex Band Structure

Christopher J. DeGrendele<sup>1,3</sup>, Jonathan A. Kazakov<sup>2,3</sup>, and Matthew G. Reuter<sup>2,3</sup>

<sup>1</sup>Department of Physics and Astronomy, <sup>2</sup>Department of Applied Mathematics & Statistics,

<sup>3</sup>Institute for Advanced Computational Science, Stony Brook University, Stony Brook, New York 11794

## Abstract and Introduction

The band structure of a material can provide many details into the properties of perfect crystalline systems, such as electrical conductivity, which has led to many real world applications, two examples are "next generation" photovoltaics and battery electrodes. However, by allowing the wavevector to be complex, we can analyze non-pristine materials where conventional band structure cannot be applied. In this study we simulate complex band structure of systems with interesting chemical properties to better understand the fundamentals of complex band structure. Through these simulations we have found interesting qualitative data that show quantum interference, asymmetric and vertical bands. Using these results can help contribute to advancements in material science and the material genome initiative by providing the mathematical language to describe these model systems.

- At the core of conventional band structure we assume translational symmetry, to which we can evoke Bloch's theorem: (1)

$$\psi_{n,\vec{k}}(\vec{r}) = e^{i\vec{k}\cdot\vec{r}} u_{n,\vec{k}}(\vec{r})$$

Wavefunction  $\psi_{n,\vec{k}}(\vec{r})$  is the product of a plane wave  $e^{i\vec{k}\cdot\vec{r}}$  and a periodic function  $u_{n,\vec{k}}(\vec{r})$  of the lattice. The wave vector  $\vec{k}$  is shown as a downward arrow, and the periodic function is shown as an upward arrow.

- Physically this assumes there is infinite symmetry in all directions and the wavefunction decays exponentially as we travel from cell to cell (when outside of a band).

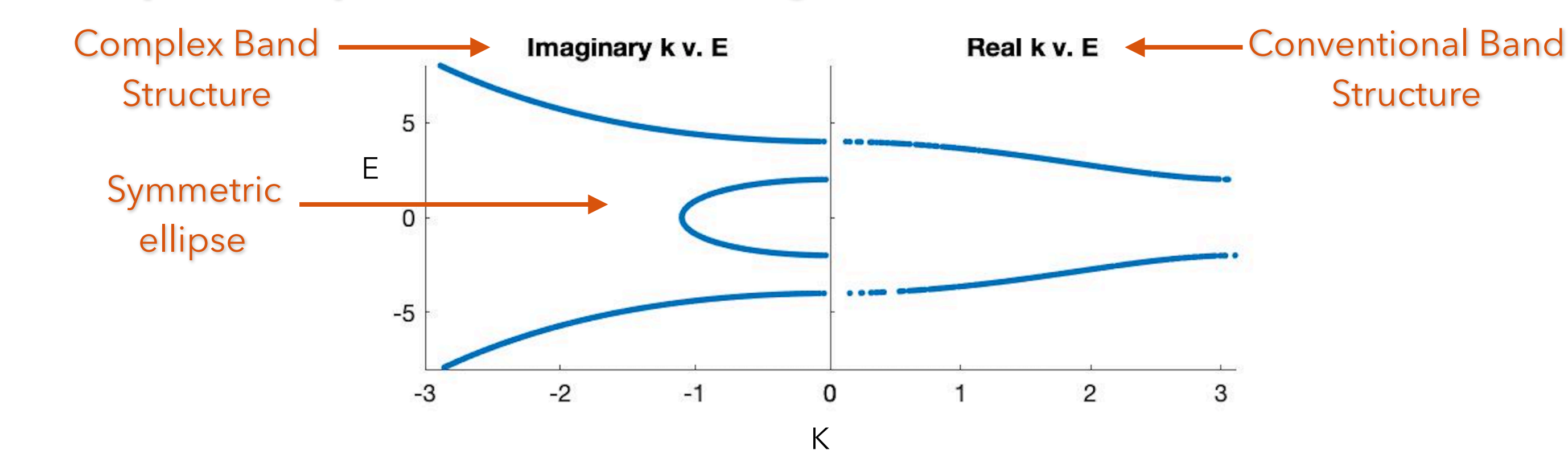
Coupling to next layer

$$\mathbf{H} = \begin{bmatrix} \mathbf{H}_D & \mathbf{H}_S^\dagger & 0 & 0 & \dots \\ \mathbf{H}_S & \mathbf{H}_D & \mathbf{H}_S^\dagger & 0 & \dots \\ 0 & \mathbf{H}_S & \mathbf{H}_D & \mathbf{H}_S^\dagger & \dots \\ \vdots & \vdots & \vdots & \vdots & \ddots \end{bmatrix}$$

Current layer  $\mathbf{H}_D$  and  $\mathbf{H}_S$  are shown as arrows pointing to the corresponding blocks in the matrix. A molecular structure diagram is shown to the right.

- The coupling between cells and the cells themselves becomes identical, represented as a hermitian and block tridiagonal matrix.
- Mathematically, these are easier to work with
- When we break translational symmetry, such as considering a finite version, the matrix is no longer infinitely long and we are interested in the physical effects of this.

## A Simple Example (Su-Schrieffer-Heeger Model) (4)



## Simulation and Methods

We can analyze the band structure of a material by making use of the transfer matrix. (1)

$$\mathbf{T} = \begin{bmatrix} \mathbf{H}_S^\dagger(\mathbf{E}\mathbf{I} - \mathbf{H}_D) & -\mathbf{H}_S^\dagger\mathbf{H}_S \\ \mathbf{I} & \mathbf{0} \end{bmatrix}$$

- $\mathbf{T}(E)$  connects wavefunctions from one layer to the next.
- The complex band structure comes from the eigenvalues, or in the non symmetric case, the spectra, of this  $\mathbf{T}(E)$ .

### Key Questions

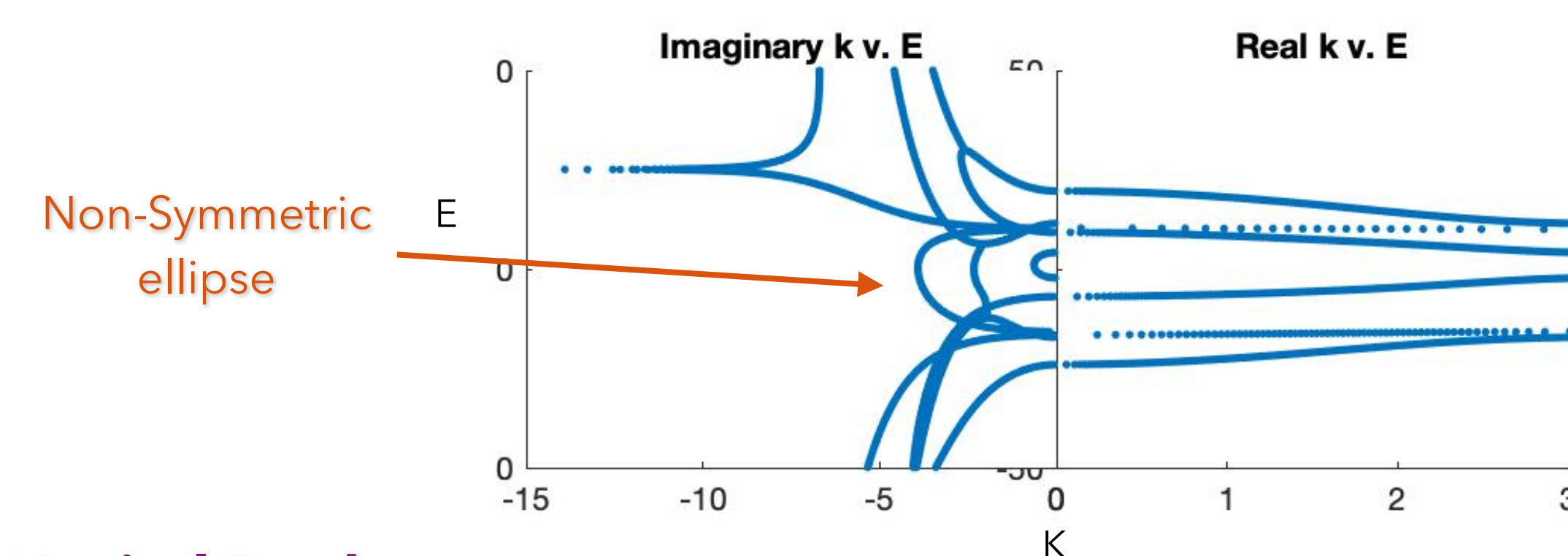
- Can CBS provide additional insight into observed physical phenomena?
- How does  $\|\mathbf{T}(E)\|$  evolve with  $k$ ?

- We can simplify the problem greatly if we note that the diagonal matrix  $\mathbf{H}_S$  represents the coupling to the other cells.
  - Taking the rank( $\mathbf{H}_S$ ) we can delete all the eigenvalues that do not effect this structure.

## Characterizing Behaviors

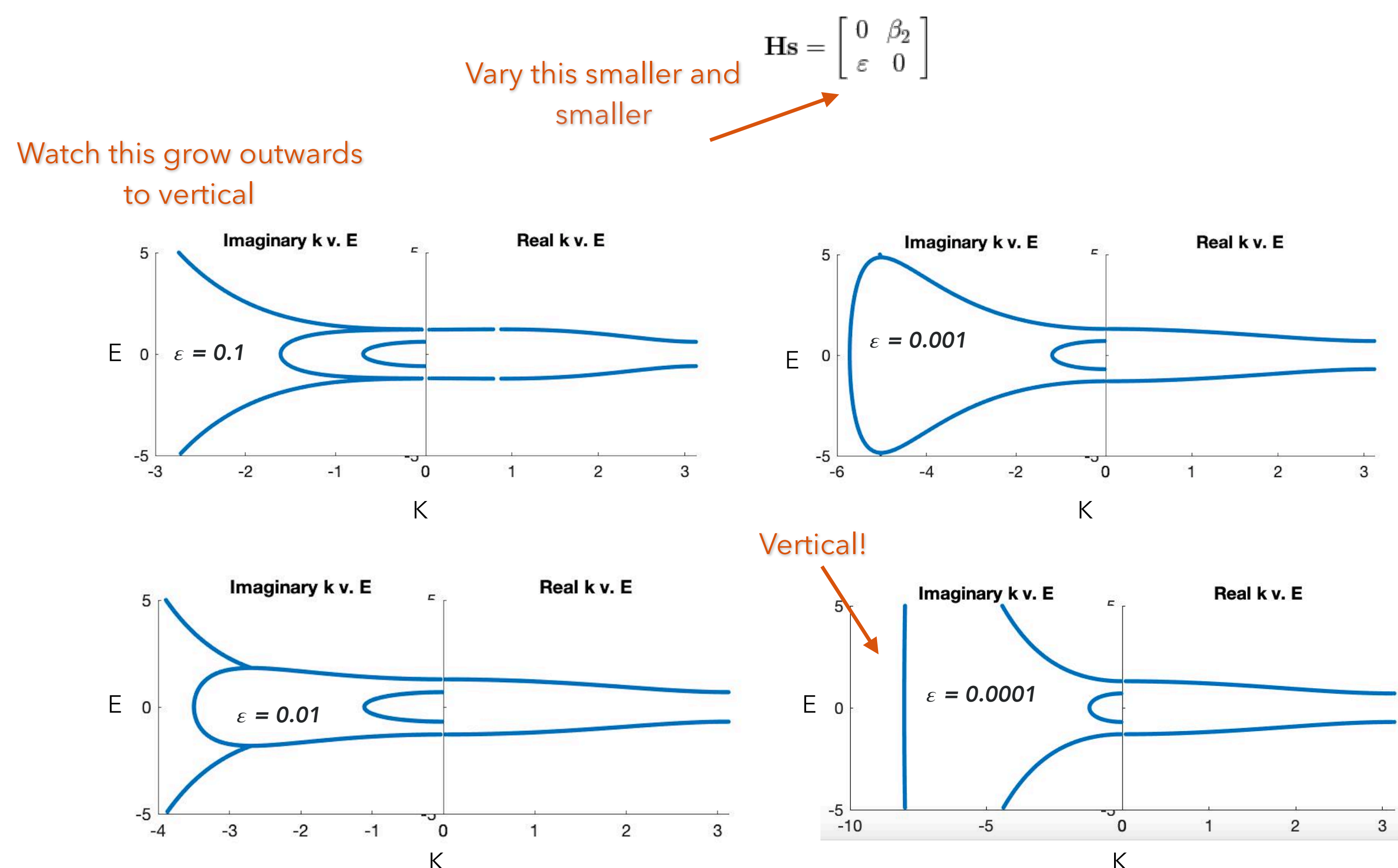
### Asymmetric Complex Bands

- More sophisticated computational studies have reported asymmetric complex bands
- A tight binding model of siloxane motivated by (2) shows asymmetric bands.
- Preliminary hypothesis: Need higher rank  $\mathbf{H}_S$  matrices



### Vertical Bands

Vertical complex bands (5) are a computational artifact and are not physical.

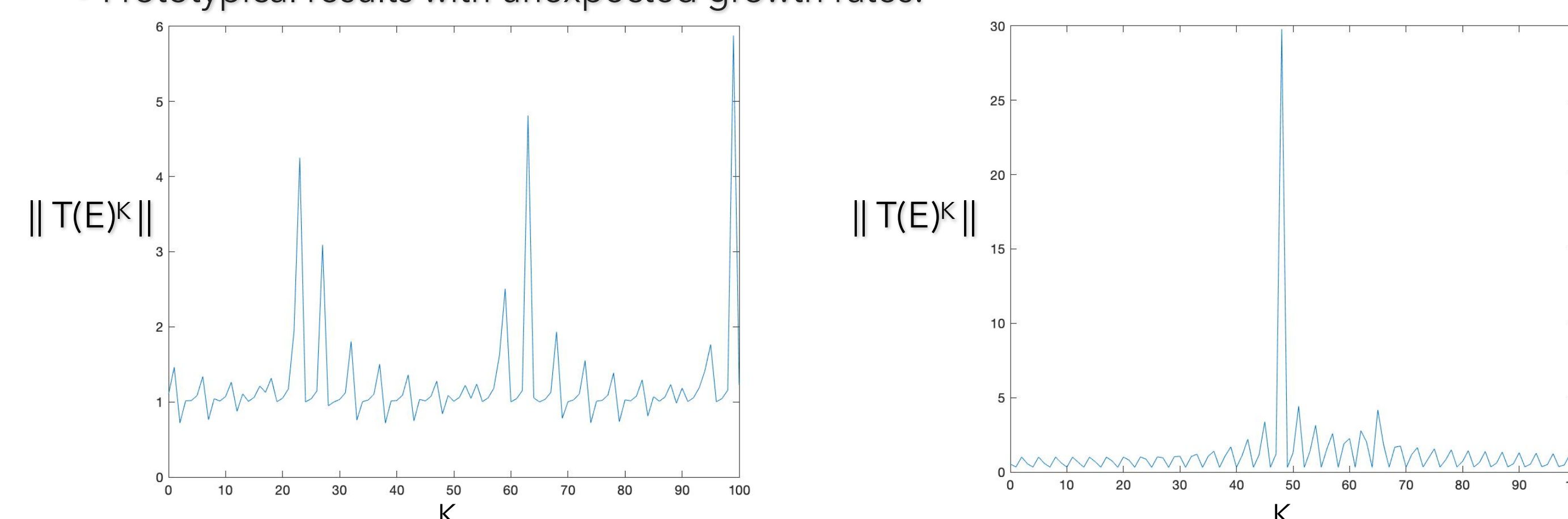


- Conclusion: We see the vertical band becoming more and more separate and apparent as we make  $\epsilon$  smaller.  $\epsilon=0$  is the true answer, the others are approximations.
- Here we found that they arise from having a singular matrix  $\mathbf{H}_S$

## Characterizing Observed Phenomena

### Transient Effects

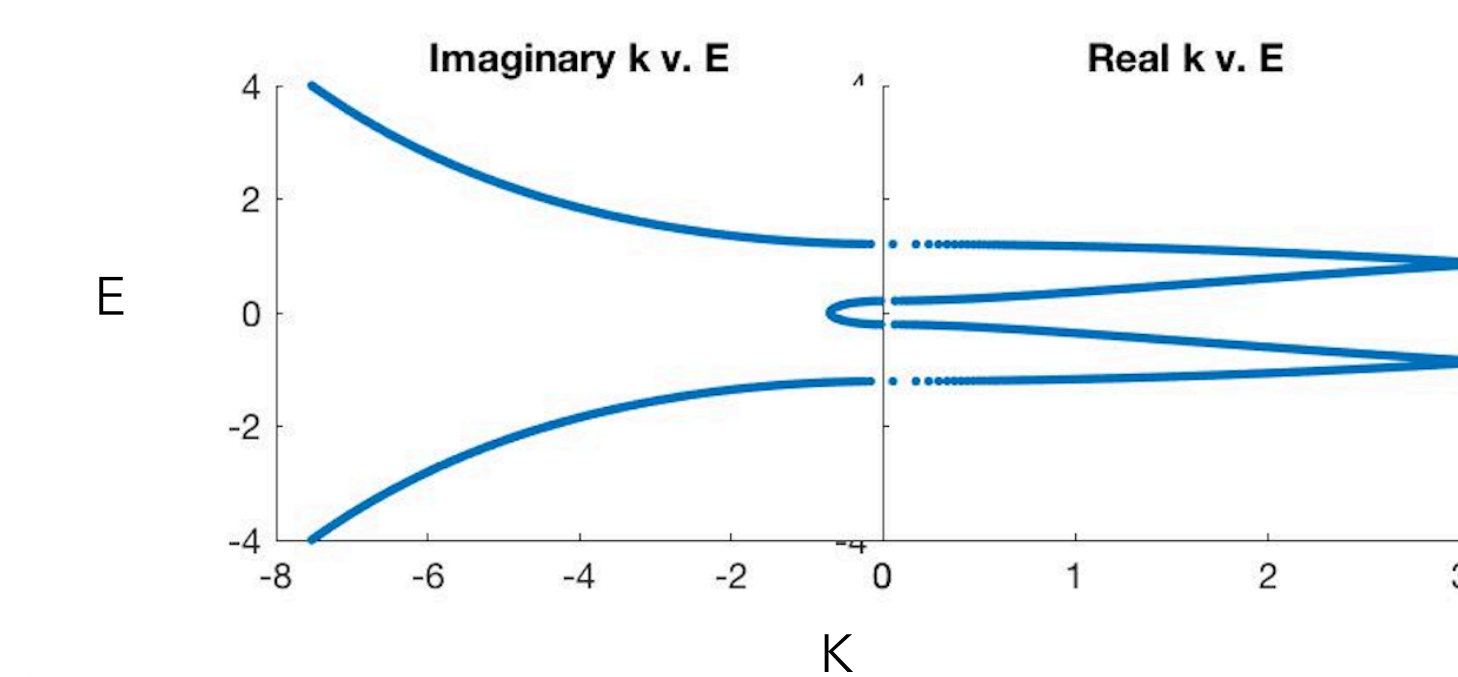
- Conventional analyses expect the wavefunction to decay exponentially between layers of the crystallized structure when outside the band
- This is not always true as observed in the electrical conductance of various polymeric molecules, which display transient effects
- Hypothesis: The unexpected increase in conductance is not caused by spurious changes in the molecular geometry, but are a fundamental consequence of the molecules' band structures.
- Prototypical results with unexpected growth rates:



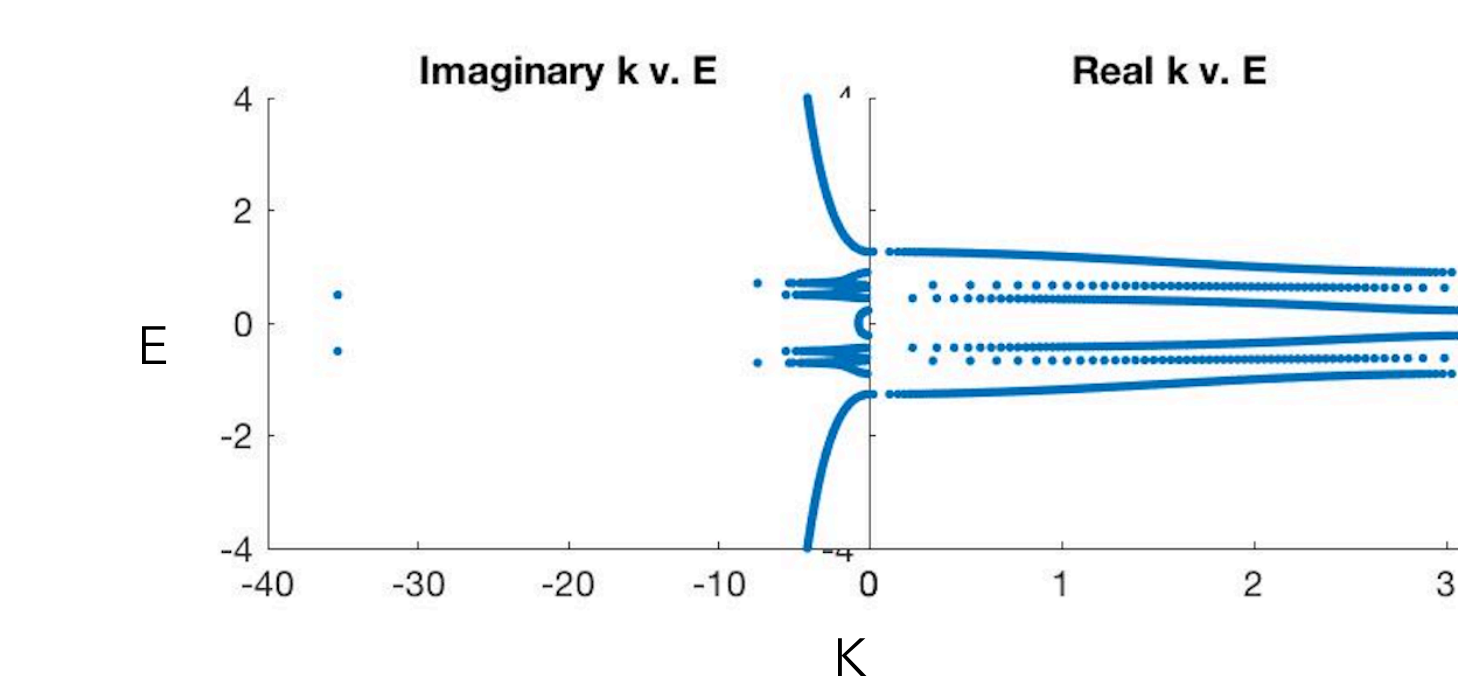
### Destructive Quantum Interference

- Destructive Quantum Interference suppresses connections in the wavefunction from one layer to another. How does it manifest in the CBS?

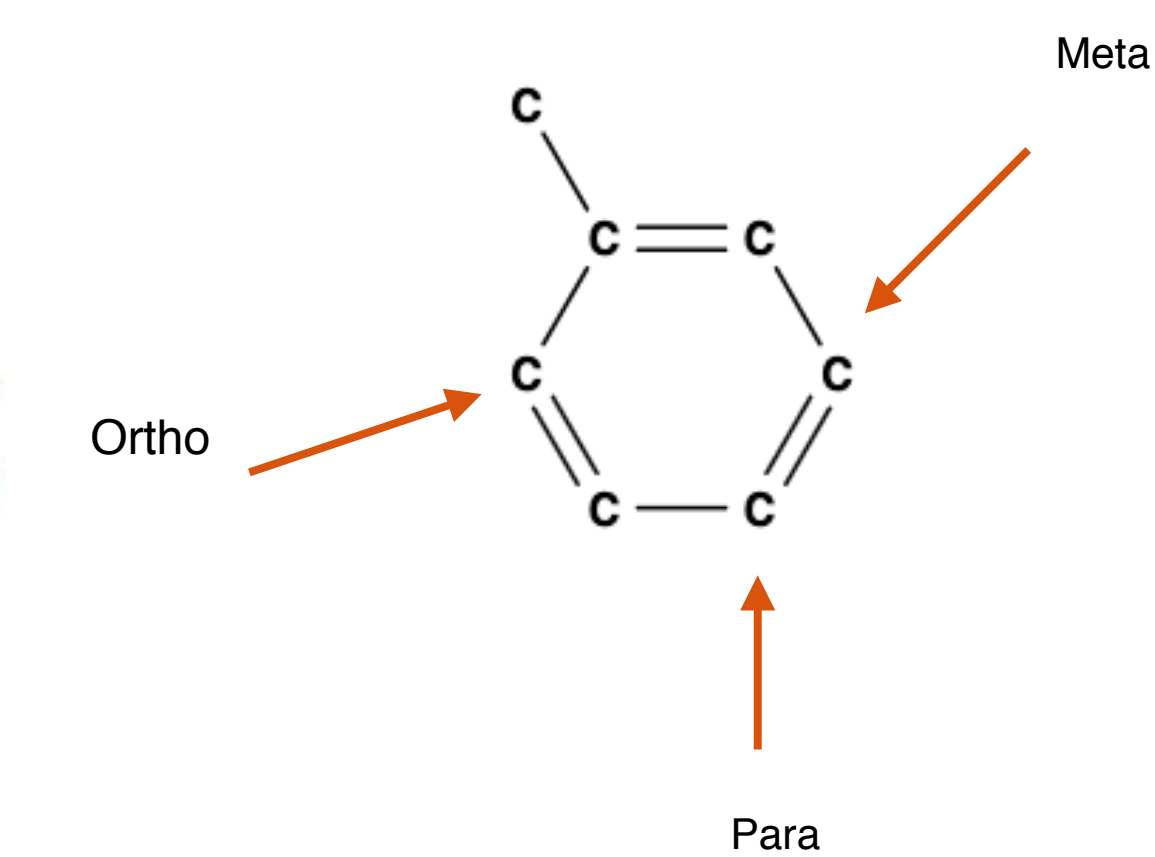
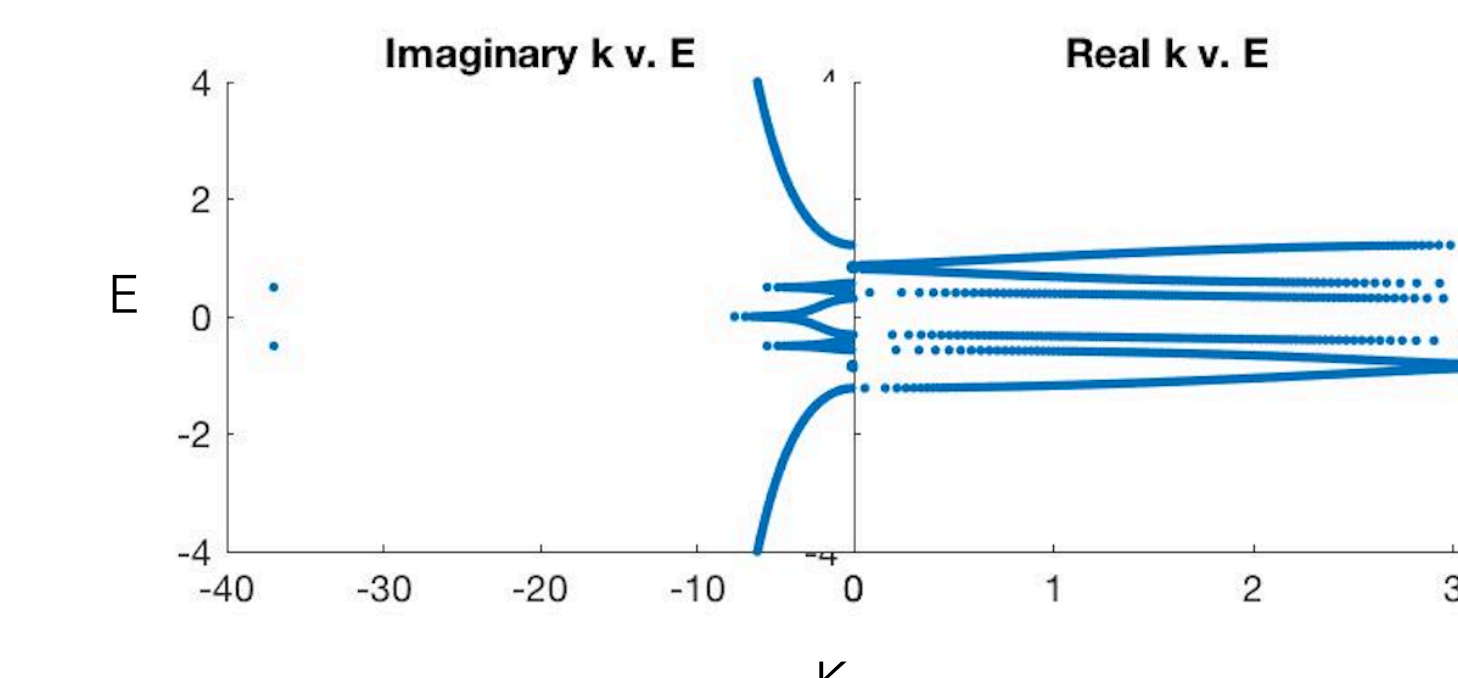
Para-benzene (3)



Ortho-benzene



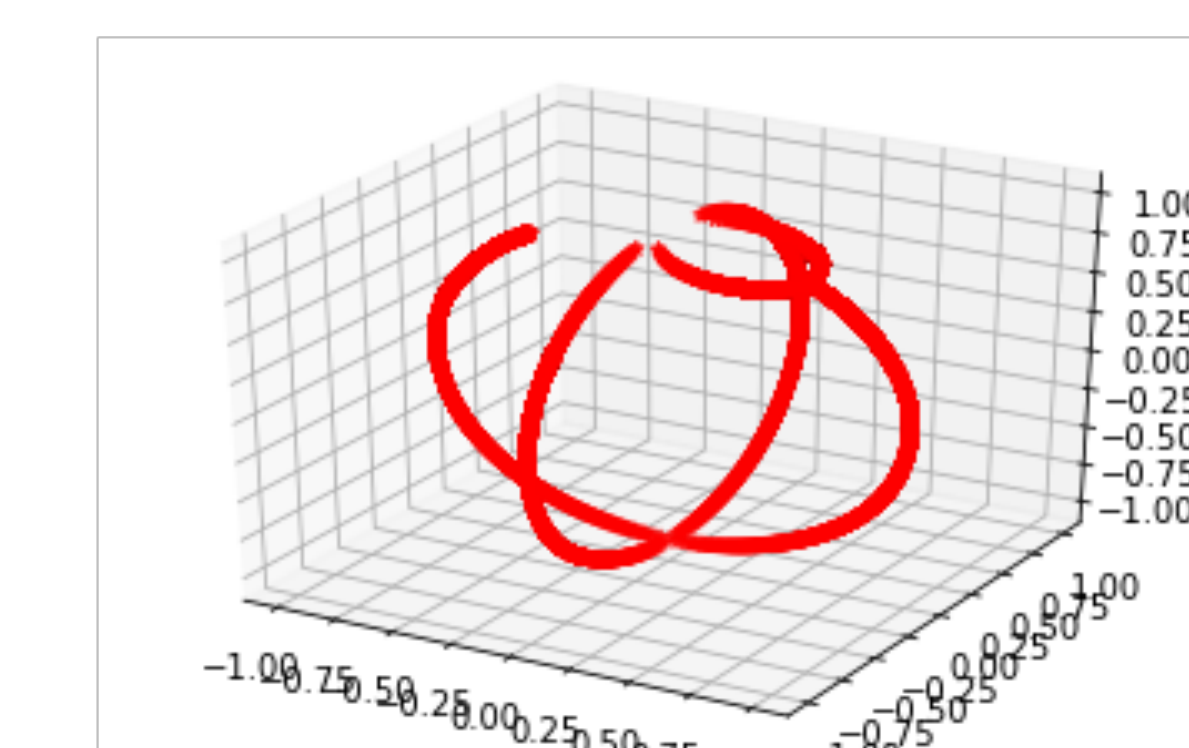
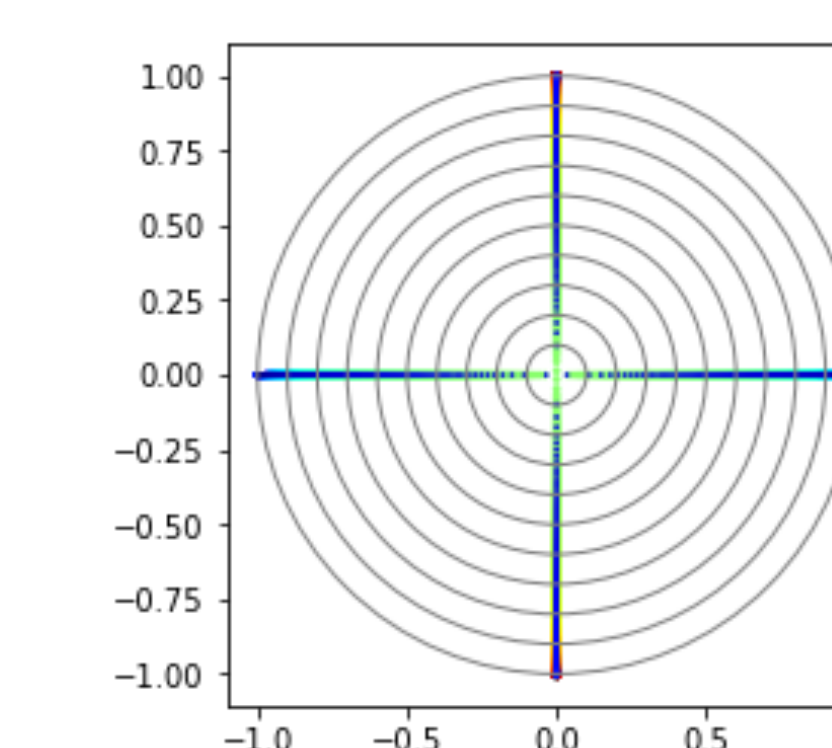
Meta-benzene



- Conclusion: Destructive Quantum Interference corresponds to infinite wavevectors,  $k = \infty$

### Stereographic Projections

- Another way of looking at destructive quantum interference is using stereographic projections.
- Project the complex data onto a Riemann sphere allows us to physically see when  $k$  crosses infinity
- Here is the Ortho-benzene molecule again:



## Conclusions

- Complex band structure allows us to describe transient effects and destructive quantum interference among other important visual behaviors
- We were able to describe numerical artifacts such as vertical bands
- Furthermore a future study could be to actually visualize our results in a laboratory setting to test whether transient effects could be seen in these chemical compounds

## References & Acknowledgments

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